

## Atomic Physics' starting point: the old Bohr model



Niels Bohr  
Nobel prize 1922

Solution to radiative instability of the atom:

- atom exists in a discrete set of stationary states
- radiative transitions → quantum jumps between levels

$$h\nu = \frac{hc}{\lambda} = E_i - E_f$$

- angular momentum is quantized

$$L = \mu v r = n\hbar$$

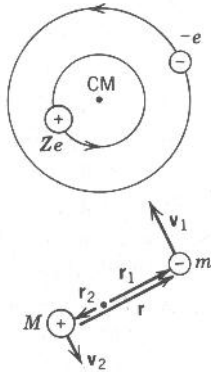
Quantized level energies

$$E_n = -\frac{Z^2}{n^2} R_H \quad \text{with} \quad R_H = \frac{\mu}{m_e} R_\infty$$



Reduced mass in two-particle problem:  $\mu = \frac{m_e M}{m_e + M}$

## Reduced mass in the old Bohr model → the one particle problem



Position vectors:

$$\vec{r}_1 = \frac{M}{m+M} \vec{r}$$

$$\vec{r}_2 = -\frac{m}{m+M} \vec{r}$$

Kinetic energy

$$K = \frac{1}{2} m_1 v_1^2 + \frac{1}{2} m_2 v_2^2 = \frac{1}{2} \mu v^2$$

Angular momentum

$$L = m_1 v_1 r_1 + m_2 v_2 r_2 = \mu v r$$

Relative coordinates:

$$\vec{r} = \vec{r}_1 - \vec{r}_2$$

Centre of Mass

$$m\vec{r}_1 + M\vec{r}_2 = 0$$

Velocity vectors:

$$\vec{v}_1 = \frac{M}{m+M} \vec{v}$$

$$\vec{v}_2 = -\frac{M}{m+M} \vec{v}$$

With reduced mass

$$\mu = \frac{mM}{m+M}$$

Centripetal force

$$F = \frac{m_1 v_1^2}{r_1} = \frac{m_2 v_2^2}{r_2} = \frac{\mu v^2}{r}$$

Relative velocity

$$\vec{v} = \frac{d\vec{r}}{dt}$$



## The old Bohr model → quantisation of radial motion

Kinetic energy  $K = \frac{1}{2} \mu v^2$

Potential energy  $V = -\frac{Ze^2}{4\pi\epsilon_0 r}$

Total energy

$$E = \frac{1}{2} \mu v^2 - \frac{Ze^2}{4\pi\epsilon_0 r}$$

System is bound by the Coulomb force

$$F_{\text{Centripetal}} = F_{\text{Coulomb}}$$

$$\frac{\mu v^2}{r} = \frac{Ze^2}{4\pi\epsilon_0 r^2}$$

Total energy (bound system)

$$E = \left( \frac{1}{2} - 1 \right) \frac{Ze^2}{4\pi\epsilon_0 r} = -\frac{1}{2} \frac{Ze^2}{4\pi\epsilon_0 r}$$

Quantisation of angular momentum:

$$L = \mu v r = n \frac{h}{2\pi} = n\hbar$$

$$v^2 r^2 = \frac{Ze^2}{4\pi\epsilon_0} \frac{r}{\mu} = \frac{n^2 \hbar^2}{\mu^2}$$

Quantisation of radius in orbit:

$$r_n = \frac{n^2}{Z} \frac{4\pi\epsilon_0 \hbar^2}{e^2 \mu} = \frac{n^2}{Z} \frac{m_e}{\mu} a_0$$

With Bohr radius  $a_0$



## The old Bohr model → quantisation of energy

$$E = -\frac{1}{2} \frac{Ze^2}{4\pi\epsilon_0 r_n}$$

$$r_n = \frac{n^2}{Z} \frac{4\pi\epsilon_0 \hbar^2}{e^2 \mu} = \frac{n^2}{Z} \frac{m_e}{\mu} a_0$$

$$E_n = -\frac{Z^2}{n^2} \left( \frac{e^2}{4\pi\epsilon_0} \right)^2 \frac{\mu}{2\hbar^2}$$

$$E_n = -\frac{Z^2}{n^2} \frac{\mu}{m_e} R_\infty$$

Rydberg constant

$$R_\infty = \left( \frac{e^2}{4\pi\epsilon_0} \right)^2 \frac{m_e}{2\hbar^2} = E_0$$



Energy levels in the Bohr model:

$$E_n = -\frac{Z^2}{n^2} \frac{\mu}{m_e} R_\infty$$

Rewrite:

$$E_n = -\frac{Z^2}{n^2} \left( \frac{\mu}{m_e} \right) \frac{1}{2} \alpha^2 m_e c^2$$

↔ ↔  
dimensionless energy

$$\alpha \approx \frac{1}{137}$$



## The old Bohr model → energy levels

$$E_n = -\frac{Z^2}{n^2} \left( \frac{\mu}{m_e} \right) \frac{1}{2} \alpha^2 m_e c^2$$

$\longleftrightarrow$   
 dimensionless energy

Energy scale in the molecule

$$R_\infty = \frac{1}{2} \alpha^2 m_e c^2$$

13.60 eV

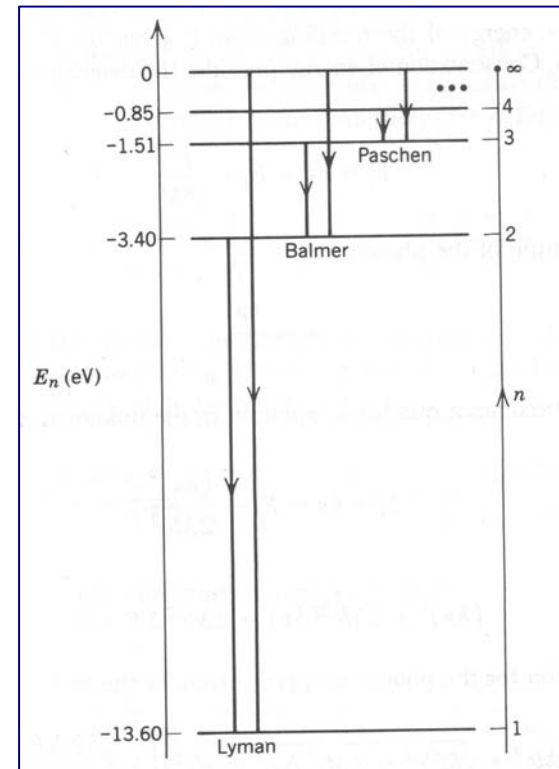
109737.31534 cm<sup>-1</sup>

Atomic Rydberg constant

$$R_H = \left( \frac{\mu}{m_e} \right) R_\infty$$



## Energy levels and Bohr "quantum jumps"

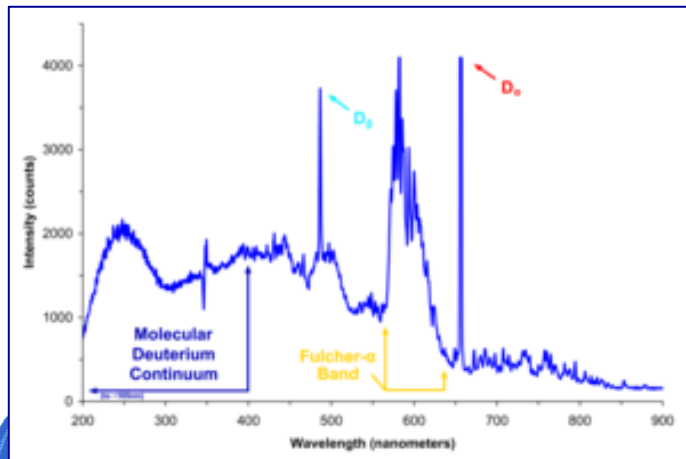


## The old Bohr model : explains Balmers and Rydberg laws



Johann Jakob Balmer

$$\lambda = \frac{h m^2}{m^2 - n^2}$$



Janne Rydberg

$$\frac{1}{\lambda} = R_H \left( \frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

Success of the Bohr model  
restricted to atomic hydrogen

Extension by Sommerfeld

## Generalisation: isotope shifts and exotic atoms

$$E_n = -\frac{Z^2}{n^2} \left( \frac{\mu}{m_e} \right) \frac{1}{2} \alpha^2 m_e c^2$$

For all atoms:  $\mu \approx m_e$

Exotic atoms:

Muonium  $(\mu^+ e^-)$

Positronium  $(e^+ e^-)$

Protonium  $(p^+ \bar{p}^-)$

Ion-pair states in molecules?  $(H^+ H^-)$

$$R_H = \left( \frac{\mu}{m_e} \right) R_\infty$$

$$\mu = \frac{mM}{m + M}$$

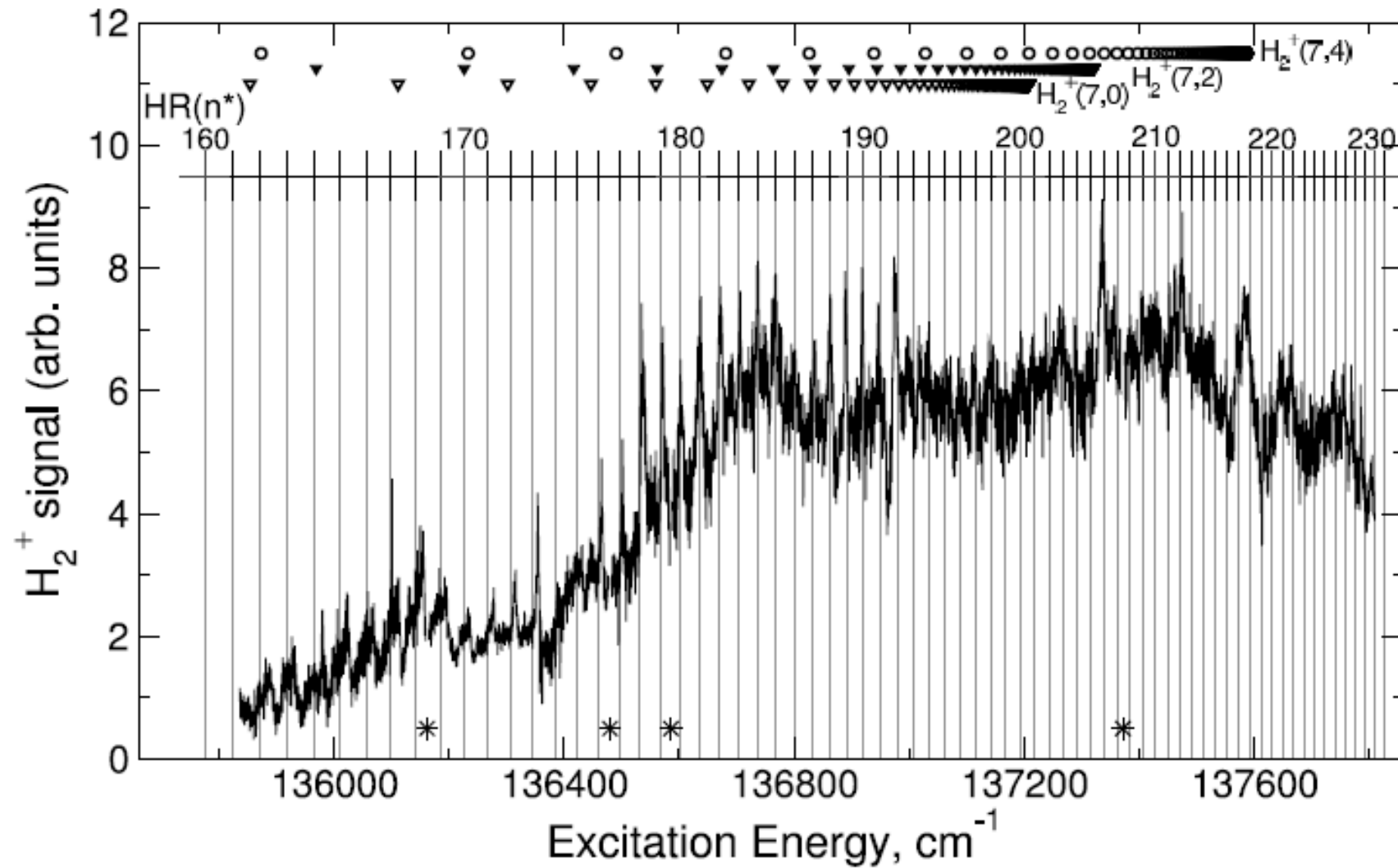


The isotope shift



The heavy Bohr atom:  $\text{H}^+\text{H}^-$

$$R_H = \left( \frac{\mu}{m_e} \right) R_\infty = 918 R_\infty$$





## Generalisation: the gravitational perspective

Gravitational potential

$$V = -G \frac{M_1 M_2}{r}$$

Coorespondence

$$\frac{Ze^2}{4\pi\epsilon_0} \leftrightarrow GM_1 M_2$$

Total energy

$$E_g = -\frac{1}{2} GM_1 M_2 \frac{1}{r}$$

Quantized radius for a gravitational system

$$r_n^g = \frac{n^2 \hbar^2}{\mu} \frac{1}{GM_1 M_2} = n^2 \left( \frac{M_1}{\mu} \right) a_g$$

Quantized energies for a gravitational system

$$E_n = -\frac{1}{n^2} \left( \frac{\mu}{M_1} \right) R_\infty^g$$

Gravitational Rydberg constant

$$R_\infty^g = (GM_1 M_2)^2 \frac{M_1}{2\hbar^2}$$



## Gravitation, example: Sun-Earth system

### Parameters

$$M_1 = 5.98 \times 10^{24} \text{ kg}$$

$$M_2 = 1.989 \times 10^{30} \text{ kg}$$

$$R_{\text{orbit}} = 150 \times 10^9 \text{ m}$$

$$G = 6.673 \times 10^{-11} \frac{\text{m}^3}{\text{s}^2 \text{ kg}}$$

$$\hbar = 1.055 \times 10^{-34} \frac{\text{kg m}^2}{\text{s}}$$

### Calculate

$$GM_1 M_2 = 7.9 \times 10^{44} \frac{\text{kg m}^3}{\text{s}^2}$$

### Bohr radius

$$a_g = 2 \times 10^{-138} \text{ m}$$

### Principal quantum number for given orbit:

$$n = 2.7 \times 10^{74}$$

### Rydberg constant

$$R_{\infty}^g = 1.7 \times 10^{182} \text{ J}$$

### Energy quanta

$$\Delta E_n = \frac{2}{n^3} R_{\infty}^g = 2 \times 10^{-41} \text{ J}$$



## Gravitation, further

Calculate the Rydberg constant  
for attraction between  
two hydrogen atoms

$$R_{eff}^g = 2.62 \times 10^{-87} J = 6.60 \times 10^{-65} cm^{-1}$$

"Lyman-alpha" transition:  $n=1 \rightarrow n=2$

$$\nu = 1.48 \times 10^{-54} Hz$$

"Bohr radius"

$$a_g^H = 3.5 \times 10^{22} m = 3.7 \times 10^7 ly$$

How about a neutron in the  
field of the Earth:  
Space quantisation possible ?

Nesvizhevsky et al

Nature 415, 297 (2002)

