

Halving the Casimir force with Conductive Oxides

S. de Man, K. Heeck, R. J. Wijngaarden, and D. Iannuzzi*

Department of Physics and Astronomy, VU University Amsterdam, De Boelelaan 1081, 1081 HV Amsterdam, The Netherlands
(Received 23 February 2009; published 22 July 2009)

The possibility to modify the strength of the Casimir effect by tailoring the dielectric functions of the interacting surfaces is regarded as a unique opportunity in the development of micro- and nanoelectromechanical systems. In air, however, one expects that, unless noble metals are used, the electrostatic force arising from trapped charges overcomes the Casimir attraction, leaving no room for exploitation of Casimir force engineering at ambient conditions. Here we show that, in the presence of a conductive oxide, the Casimir force can be the dominant interaction even in air, and that the use of conductive oxides allows one to reduce the Casimir force up to a factor of 2 when compared to noble metals.

DOI: 10.1103/PhysRevLett.103.040402

PACS numbers: 03.65.-w

The mechanical parts of micro- and nanoelectromechanical systems (MEMS and NEMS) are often designed to work at separations where the Casimir effect [1] might play a relevant role [2–5]. It is thus commonly believed that, if one could suitably engineer the strength of the Casimir force, unprecedented opportunities would come available for the development of conceptually new MEMS and NEMS [6–9]. The most simple approach to tailor the Casimir force is to properly choose the materials of which the interacting surfaces are made. According to the Lifshitz theory [10], the interaction between two objects depends on their dielectric functions. Transparent dielectrics, for example, attract less than reflective mirrors. This property may be used to reduce the Casimir attraction whenever the design requires a smaller short range interaction. It is, however, fair to say that, for the vast majority of applications, MEMS and NEMS operate in air, where surfaces tend to accumulate trapped charges. Those charges give rise to a strong electrostatic interaction that cannot be compensated by a counterbias voltage and that typically overcomes the Casimir force. It is thus difficult to imagine that the Casimir force can play an important role in MEMS and NEMS operating in air, unless all surfaces are coated with noble metals to reduce the forces due to surface charges to negligible levels. In that case, however, there is not much room to tune the strength of the Casimir interaction because the diversity in the dielectric functions of different metals is simply not large enough [11–15]. As a matter of fact, to date, there is no experiment that shows that, in air, the Casimir force can still be tuned significantly while remaining the dominant interaction mechanism.

In this Letter we present a precise measurement of the Casimir force between a gold coated sphere and a glass plate coated with either a thick gold layer or a highly conductive, transparent oxide film. The experiment was performed in air, and no electrostatic force due to residual trapped charges was observed over several weeks of measurements in either case. The decrease of the Casimir force due to the different dielectric properties of the reflective

gold layer and the transparent oxide film resulted to be as high as $\approx 40\%$ – 50% at all separations (from ≈ 50 to ≈ 150 nm). Our experiment shows that, in the presence of a conductive oxide layer, the Casimir force can still be the dominant interaction mechanism even in air, and indicates that, whenever the design might require it, it is possible to tune the Casimir attraction by a factor of 2.

Our experimental setup is designed to perform precise measurements of surface forces between a $100\ \mu\text{m}$ radius sphere and a plate as a function of their separation (see Fig. 1 and [16]). The sphere is glued onto the hanging end of a micromachined cantilever (spring constant $\approx 1\ \text{N/m}$, resonance frequency (with the sphere attached) $\approx 1.9\ \text{kHz}$). The plate is mounted on a capacitive feedback controlled piezoelectric transducer that allows one to ac-

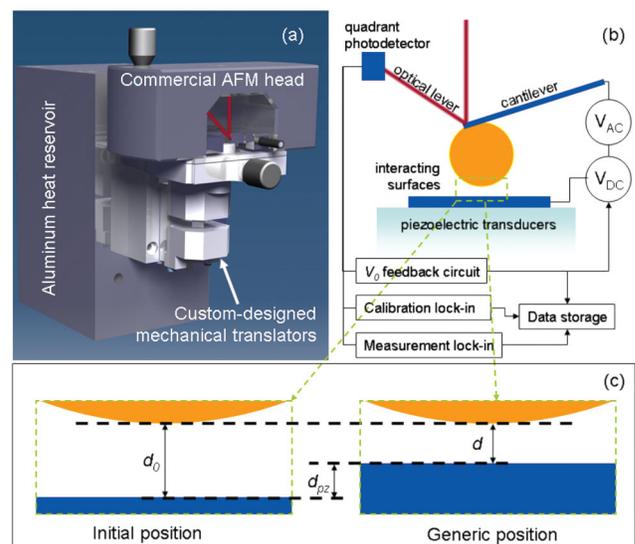


FIG. 1 (color). (a) Drawing of the experimental setup. (b) Schematic representation of the working principle of the experimental technique. (c) Definition of d_0 (initial separation), d_{pz} (movement of the piezoelectric stage), and d (separation between the two surfaces).

curately vary the separation between the sphere and the plate in discrete steps. Any force acting between the two surfaces results in a bending of the cantilever that is detected by the optical lever of a commercial atomic force microscope (AFM) head [17]. The setup is kept at a fixed temperature to reduce mechanical drifts and is placed on an active antivibration stage inside an acoustic isolation box to decouple the force sensor from external vibrations.

Precise Casimir force measurements typically require careful analysis of three crucial issues. First, even when both the sphere and the plate are coated with metallic films, there might still exist an electrostatic potential difference V_0 between the two surfaces that gives rise to a residual electrostatic force. This force must be actively canceled by counter biasing V_0 with an externally applied voltage. Second, although the relative displacements of the piezoelectric transducer that moves the plate, d_{pz} , are precisely controlled, the separation between the two interacting surfaces at the start of the measurement, d_0 , is *a priori* unknown (see Fig. 1). Therefore, the absolute separation $d = d_0 - d_{pz}$ has to be obtained from a calibration procedure. Third, the electronic signal coming out of the AFM head must be converted into force. It is thus necessary to calibrate the instrument with a controlled force. To address these problems, we have designed a measurement technique that allows one to simultaneously: (i) compensate for the residual potential, (ii) calibrate the setup, and (iii) measure the Casimir force [16,18]. In a nutshell, in each calibration/measurement run the plate is moved in discrete steps from d_0 to a minimum value of d (just before contact with the sphere). A calibrated ac electrostatic potential is applied between the sphere and the plate at a frequency ω_1 much lower than the resonance frequency of the force sensor ($\frac{\omega_1}{2\pi} = 72.2$ Hz). This ac excitation is used to drive a feedback circuit that compensates for the residual voltage V_0 , and, simultaneously, generates an electrostatic force that makes the cantilever oscillate at $2\omega_1$. The amplitude of the oscillations, measured with a lock-in amplifier [calibration lock-in in Fig. 1(b)], are recorded as a function of d_{pz} , and are then used to calibrate the instrument and extract d_0 . At the same time, a transducer mechanically coupled to the piezoelectric translator makes the plate move around d_{pz} with an amplitude of 3.85 ± 0.08 nm at a frequency ω_2 , which is again much lower than the resonance frequency of the force sensor ($\frac{\omega_2}{2\pi} = 119$ Hz) [19]. In the presence of a force that depends on separation (e.g., the Casimir force), the cantilever bends in phase with the modulation of d . The amplitude of the in-phase oscillation, measured with another lock-in amplifier [measurement lock-in in Fig. 1(b)], is proportional to the derivative of the force with respect to d . Furthermore, the presence of the cushion of air between the two surfaces gives rise to a hydrodynamic force that depends on the velocity with which the plate moves. The signal produced by this force manifests itself at the same frequency at

which d is modulated, but with a phase rotated by 90° . This contribution does not influence the output of the in-phase component and can be measured independently with the same lock-in amplifier. The integration times of the lock-in amplifiers are 8 s for every value of d_{pz} . A typical measurement run consists of ≈ 50 d_{pz} set-points in the measurement range $50 < d < 1100$ nm, and takes roughly 7 min. The cantilever responses to the modulations at ω_1 and ω_2 are < 80 pm (root-mean-square) during the entire experiment. All force measurements are performed in air at atmospheric pressure, temperature 300 K, and 29% relative humidity.

In this Letter we present two experiments performed with the same sphere and two different plates. The sphere is a polystyrene sphere with nominal radius $100 \mu\text{m}$ coated with a Ti adhesion layer followed by a 100 nm Au film (surface roughness 3.8 nm RMS). The plate used in the first experiment is a polished sapphire substrate coated with a metallic film similar to the one deposited on the sphere (surface roughness 0.8 nm RMS). The plate used in the second experiment is a float glass substrate with a 190 nm indium tin oxide (ITO, $\text{In}_2\text{O}_3:\text{Sn}$) sputtered thin film on top (PGO CEC010S, typically $8.5 \Omega/\square$, or, equivalently, $\rho = 1.6 \times 10^{-4} \Omega\text{cm}$, total surface roughness 4 nm RMS). We have measured the reflection and transmission spectra of both plates in the wavelength range

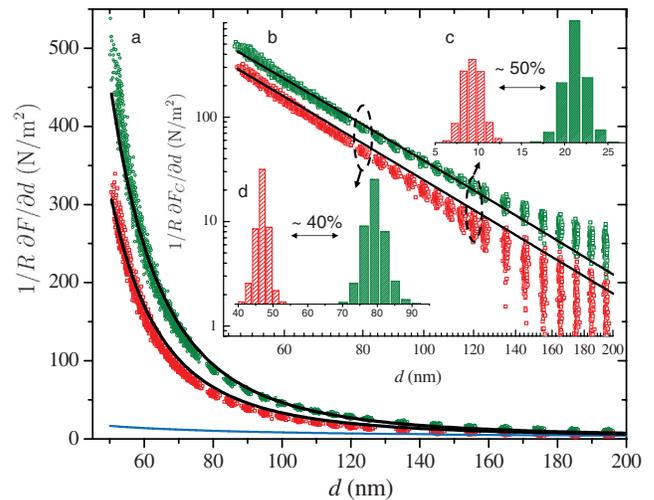


FIG. 2 (color). (a) Spatial derivative of the total force as a function of absolute surface separation for the Au-Au (green squares) and Au-ITO (red squares) interactions for randomly chosen subsets of the data (150 out of 580 for both cases). The blue line represents the derivative of the electrostatic force caused by the simultaneous calibration procedure (common to both the gold and ITO measurements). The black lines indicate the calculated Casimir forces with the electrostatic background added. (b) Spatial derivative of the Casimir force, with the electrostatic background subtracted from the data. The black lines correspond to the calculations of the Casimir force. (c) Histograms of 580 force measurements for Au-Au and Au-ITO at $d = 120$ nm. (d) Same as c, but for $d = 80$ nm.

180 nm $< \lambda < 2.5 \mu\text{m}$, and observed that the optical properties of our films are in agreement with the literature [20,21].

Figure 2(a) presents measurements of the force between the sphere and the plates coated with either Au or ITO [22]. The experimental data represent the spatial derivative of the total force (normalized by the sphere radius R), which is the sum of the Casimir interaction, a Coulomb interaction induced by the presence of trapped charges (if any), and an electrostatic attraction due to the ac calibration potential. The strength of the latter can be estimated from the simultaneous calibration procedure [18]. From Fig. 2, it is evident that this electrostatic contribution, which is anyway equal in both experiments (within 2%), is small compared to the total force signal. The black lines in Fig. 2(a) are computations of the Casimir force using the Lifshitz theory [23] with dielectric functions calculated as in [21] (for ITO) and [24] (for Au); the electrostatic force due to the calibration potential is added to the theory in order to compare with the raw data. The calculation of the Casimir force should only be considered approximate, because the dielectric functions of the samples are not known precisely [25], and no surface roughness corrections are applied [22]. Still, the agreement between the calculation and the data shows that the Casimir effect largely dominates any Coulomb interaction that would have been otherwise observed in the presence of a significant amount of trapped charges. Figure 2(b) shows the data and theory on a double logarithmic scale, where we have subtracted the electrostatic background due to the ac calibration potential using the simultaneously obtained calibration data. At small separations $d < 60$ nm, both data sets curve upwards, which might be a sign of surface roughness effects. At separations $d > 120$ nm, the experimental data for the Au-Au Casimir interaction start to deviate significantly from the theory because of an artefact caused by reflections from the optical lever light by the sample. This reflected light reaches the photodetector and causes a background signal that is not related to any force. This artefact is common to all optical lever based AFM techniques, and the related signal is typically assumed to be linear in the piezo extension and subtracted from the data accordingly [26]. Because the reflectivities of our two samples are so different, we prefer to refrain from such a procedure, and present the data as is. In Figs. 2(c) and 2(d) we show histograms of all the obtained measurements for the derivative of the Casimir force for two specific separations ($d = 120$ nm and $d = 80$ nm). The histograms at $d = 80$ nm can be described by Gaussians with a standard deviation of roughly 5%, which means that our method provides a precision in the mean measured Casimir force derivative of 0.2%. It is evident from the histograms that the spatial derivative of the Casimir force between a Au and an ITO surface is roughly $\approx 40\%$ – 50% smaller than between two Au surfaces. In our geometry (i.e., for separations much smaller than the radius of the sphere), the spatial derivative of the force is proportional to the pressure between two parallel plates [23]. We can thus conclude that the Casimir pressure that one would measure between a Au plate kept parallel to an ITO plate would be roughly $\approx 40\%$ – 50% smaller than in the case of two Au plates.

Even though the agreement between the theoretical prediction and the measurement of the Casimir force in both situations is good, one might still argue that the observed decrease could be mimicked by drifts in d_0 . For both measurements series, we gathered 580 data sets continuously, which allows us to directly assess the run-to-run drift of d_0 . Because of the temperature stabilization of our setup, the mechanical drift is very small at ≈ 0.1 nm and ≈ 0.2 nm per run for the Au-Au and the Au-ITO experiments, respectively. We conclude that the decrease of the Casimir force cannot be ascribed to drifts in d_0 . We have also verified that the electrostatic force used to calibrate the instrument and extract d_0 follows what expected from elementary electrostatic arguments, as suggested in [27] and discussed in [16]. Concerning the compensation voltage, we observed that V_0 varies approximately 1 mV and 3 mV over the complete measurement range in the Au-Au and Au-ITO cases, respectively [28]. These slight variations of V_0 do not compromise the measurement of the Casimir force at the current level of sensitivity.

Finally, the different surface roughnesses of the sphere and the plates also influence the strength of the Casimir effect. Since both experiments are conducted with the same sphere, the difference in the observed Casimir force can never be due to the surface roughness of the sphere. Second, we recall that the surface roughness of the ITO sample is larger than that of the Au substrate. Since surface roughness enhances the Casimir force [29] we note that, if it played a role in the probed separation range, it would

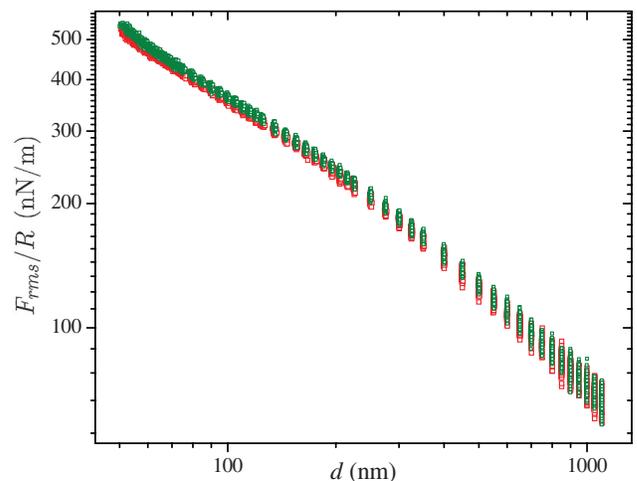


FIG. 3 (color). Hydrodynamic force acting on the sphere as a function of the absolute separation d for Au-Au (green squares) and Au-ITO (red squares), for the data sets shown in Figs. 2(a) and 2(b).

lead to a stronger interaction between Au and ITO than between two Au surfaces, contrary to the measurements presented in Fig. 2.

To make our claim even more robust, we can now compare the hydrodynamic force observed during the two experiments. Because the geometrical configuration of the experiment is equal in both cases, we expect to measure the same hydrodynamic force. The results are shown in Fig. 3. It is clear that the hydrodynamic force is very similar in both cases, although there exists a slight discrepancy between the two curves ($\approx 2\%$). Since both curves are parallel on a double logarithmic scale, we conclude that this discrepancy cannot be ascribed to a difference in the calibration of d_0 . Therefore, we rule out that the large difference in the Casimir force reported in Fig. 2 be due to an error in the determination of d_0 .

In conclusion, we have demonstrated that an ITO coating of one of the two surfaces is sufficient to readily create situations where the Casimir force is still the dominant interaction mechanism regardless the presence of air in the surroundings. Since ITO is transparent over a wide range of frequencies, the Casimir attraction is up to a factor of 2 smaller when compared to the case of the Au-Au interaction, leaving ample room for Casimir force engineering even at ambient conditions, where MEMS and NEMS typically operate.

This work was supported by the Netherlands Organisation for Scientific Research (NWO), under the Innovational Research Incentives Scheme VIDI-680-47-209. The authors thank A. Baldi, F. Mul, U. Mohideen, J. H. Rector, B. Dam, and R. Griessen for useful discussions, and the ESF Research Network CASIMIR for providing visit grants. S.M. acknowledges financial support from the Kavli Institute for Theoretical Physics (NSF PHY05-51164). D.I. acknowledges financial support from the European Research Council under the European Community's Seventh Framework Programme (FP7/2007-2013)/ERC Grant Agreement 201739.

*iannuzzi@few.vu.nl

- [1] H. B. G. Casimir, Proc. K. Ned. Akad. Wet. **51**, 793 (1948).
- [2] F. W. Delrio *et al.*, Nature Mater. **4**, 629 (2005).
- [3] H. B. Chan *et al.*, Science **291**, 1941 (2001).
- [4] H. B. Chan *et al.*, Phys. Rev. Lett. **87**, 211801 (2001).
- [5] F. M. Serry, D. Walliser, and G. J. Maclay, J. Appl. Phys. **84**, 2501 (1998).
- [6] F. Capasso *et al.*, IEEE J. Sel. Top. Quantum Electron. **13**, 400 (2007).
- [7] O. Kenneth *et al.*, Phys. Rev. Lett. **89**, 033001 (2002).
- [8] J. N. Munday *et al.*, Phys. Rev. A **71**, 042102 (2005).
- [9] J. N. Munday, F. Capasso, and V. A. Parsegian, Nature (London) **457**, 170 (2009).
- [10] E. M. Lifshitz, Sov. Phys. JETP **2**, 73 (1956).
- [11] R. S. Decca *et al.*, Phys. Rev. Lett. **91**, 050402 (2003).
- [12] An unsuccessful attempt to change the Casimir force in gas is reported in D. Iannuzzi, M. Lisanti, and F. Capasso, Proc. Natl. Acad. Sci. U.S.A. **101**, 4019 (2004).
- [13] $\approx 1\%$ tunability of the Casimir force in vacuum was shown in F. Chen *et al.*, Phys. Rev. B **76**, 035338 (2007).
- [14] $\approx 20\%$ tunability of the Casimir force was observed in M. Lisanti, D. Iannuzzi, and F. Capasso, Proc. Natl. Acad. Sci. U.S.A. **102**, 11989 (2005); That method is not practical, because the interaction strength strongly varies on the local thickness of the film (see G. Jourdan, A. Lambrecht, F. Comin, and J. Chevrier, Europhys. Lett. **85**, 31001 (2009)).
- [15] $\approx 30\%$ variations of the Casimir force were observed in J. Chen *et al.*, Phys. Rev. A **72**, 020101(R) (2005). Also this strategy is not practical for MEMS in air, because the passivation layer is not stable under atmospheric conditions.
- [16] S. de Man, K. Heeck, and D. Iannuzzi, Phys. Rev. A **79**, 024102 (2009).
- [17] Measurements of surface forces between a sphere and a plate with an AFM were first reported in W. A. Ducker, T. J. Senden, and R. M. Pashley, Nature (London) **353**, 239 (1991). Its utilization for large range forces between metals was first introduced in [26].
- [18] S. de Man *et al.* (to be published).
- [19] A similar technique was introduced in G. Bressi *et al.*, Classical Quantum Gravity **18**, 3943 (2001).
- [20] E. D. Palik, *Handbook of Optical Constants of Solids* (Academic Press, San Diego, 1998).
- [21] H. Fujiwara and M. Kondo, Phys. Rev. B **71**, 075109 (2005), with $N_{\text{Hall}} = 1.2 \times 10^{21} \text{ cm}^{-3}$.
- [22] Our estimate of d_0 relies on the *proximity force approximation* [23], which is only valid for $d \ll R$. This assumption is not entirely correct in the probed separation range [16] and results in a systematic error in d_0 of $\approx 1.4 \text{ nm}$. Still, the corresponding underestimate of the separation is equal for both Au-Au and Au-ITO, and can thus be neglected in the comparison.
- [23] V. A. Parsegian, *Van der Waals Forces* (Cambridge University Press, New York, 2006).
- [24] S. de Man and D. Iannuzzi, New J. Phys. **8**, 235 (2006).
- [25] V. B. Svetovoy *et al.*, Phys. Rev. B **77**, 035439 (2008).
- [26] U. Mohideen and A. Roy, Phys. Rev. Lett. **81**, 4549 (1998).
- [27] W. J. Kim *et al.*, Phys. Rev. A **78**, 020101(R) (2008).
- [28] The value of V_0 at $d = 100 \text{ nm}$ during the entire experiment drifts from -106 to -103 mV for Au and from 72 to 50 mV for ITO.
- [29] A. A. Maradudin and P. Mazur, Phys. Rev. B **22**, 1677 (1980).